

Nonparametric Inference for Right-Censored Data
Using Smoothing Splines

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The following regularity conditions are needed to prove the main results.

(C1) The probability $P(C \geq 1) > 0$.

(C2) The censoring time C and the survival time T are independent.

(C3) The hazard function of T , $\lambda(t)$ is bounded away from 0 and ∞ , that is, there exist constants $b_1 > 0$ and $b_2 < \infty$ such that $b_1 \leq \lambda(t) \leq b_2$.

To prove Lemma 1, we define the inner product $\langle \cdot, \cdot \rangle_1$ as a special case of $\langle \cdot, \cdot \rangle_\lambda$ when $\lambda = 1$ and the corresponding norm $\|g\|_1^2 = \langle g, g \rangle_1$.

Proof of Lemma 1. Let $g_n(t)$ be the B-spline function satisfying $\|g_n - g_0\|_\infty = O(n^{-vm})$. It follows from (P_2) , we can choose $h_n \in \Psi_{m,\mathcal{I}}$ satisfying

$\|h_n\|_1 = O(n^{-(1-v)/2} + n^{-vm})$ and $\|h_n\|_\infty = O(n^{-(1-v)/2} + n^{-vm})$. For $\alpha \in R$, write

$$\begin{aligned} H_n(\alpha) &= l_{n,\lambda}(g_n + \alpha h_n) \\ &= - \int_{\mathbb{I}} \exp\{g_n(t) + \alpha h_n(t)\} S_n(t) dt + \frac{1}{n} \sum_{i=1}^n \Delta_i(g_n + \alpha h_n)(Y_i) \\ &\quad - \frac{1}{2} \langle W_\lambda(g_n + \alpha h_n), g_n + \alpha h_n \rangle_\lambda. \end{aligned}$$

The derivative of $H_n(\alpha)$ with respect to α is

$$\begin{aligned} H'_n(\alpha) &= - \int_{\mathbb{I}} \exp\{g_n(t) + \alpha h_n(t)\} h_n(t) S_n(t) dt + \frac{1}{n} \sum_{i=1}^n \Delta_i h_n(Y_i) \\ &\quad - \alpha \langle W_\lambda h_n, h_n \rangle_\lambda - \langle W_\lambda g_n, h_n \rangle_\lambda \\ &= - \int_{\mathbb{I}} [\exp\{g_0(t)\} + \exp\{g_0(t)\} \{g_n(t) - g_0(t) + \alpha h_n(t)\} \{1 + o(1)\}] \\ &\quad h_n(t) S_n(t) dt - \alpha \langle W_\lambda h_n, h_n \rangle_\lambda - \langle W_\lambda g_n, h_n \rangle_\lambda + \frac{1}{n} \sum_{i=1}^n \Delta_i h_n(Y_i) \\ &= -\alpha \int_{\mathbb{I}} h_n^2(t) \{1 + o(1)\} \exp\{g_0(t)\} S_n(t) dt - \alpha \langle W_\lambda h_n, h_n \rangle_\lambda \\ &\quad - \left[\int_{\mathbb{I}} \exp\{g_0(t)\} h_n(t) S_n(t) dt - \frac{1}{n} \sum_{i=1}^n \Delta_i h_n(Y_i) \right] \\ &\quad - \int_{\mathbb{I}} \{g_n(t) - g_0(t)\} \{1 + o(1)\} h_n(t) \exp\{g_0(t)\} S_n(t) dt - \langle W_\lambda g_n, h_n \rangle_\lambda \\ &\equiv -\alpha I_1 - \alpha I_2 + I_3 + I_4 + I_5. \end{aligned}$$

It follows from the condition (C3) that

$$\begin{aligned}
|I_1| &= \left| \int_{\mathbb{I}} h_n^2(t) \{1 + o(1)\} \exp\{g_0(t)\} S_n(t) dt \right| \\
&\leq \{1 + o(1)\} b_2 \|h_n\|_{L_2}^2 \\
&= O_p(n^{-(1-v)} + n^{-2vm}),
\end{aligned}$$

where $\|\cdot\|_{L^2}$ is the L_2 norm.

Next, we consider I_3 . In view of the fact that $M_i(t)$ is a martingale, we have

$$\begin{aligned}
E\{|I_3|^2\} &= \frac{1}{n} E \left[\Delta_i h_n(Y_i) - \int_{\mathbb{I}} \exp\{g_0(t)\} h_n(t) S_n(t) dt \right]^2 \\
&= \frac{1}{n} E \left\{ \int_{\mathbb{I}} h_n(t) dM_i(t) \right\}^2 \\
&= \frac{1}{n} \int_{\mathbb{I}} h_n^2(t) \exp\{g_0(t)\} S(t) dt \\
&= O(n^{-(2-v)} + n^{-2vm-1}).
\end{aligned}$$

Thereby, we have $I_3 = O_p(n^{-\frac{(2-v)}{2}} + n^{-vm-1/2})$. On the other hand, from $\|g_n - g_0\|_{\infty} = O(n^{-vm})$, we get

$$\begin{aligned}
|I_4| &= \left| - \int_{\mathbb{I}} \{g_n(t) - g_0(t)\} \{1 + o(1)\} h_n(t) \exp\{g_0(t)\} S_n(t) dt \right| \\
&\leq \left| - \int_{\mathbb{I}} \{g_n(t) - g_0(t)\} \{1 + o(1)\} h_n(t) \exp\{g_0(t)\} \{S_n(t) - S(t)\} dt \right| \\
&\quad + \left| \int_{\mathbb{I}} \{g_n(t) - g_0(t)\} \{1 + o(1)\} h_n(t) \exp\{g_0(t)\} S(t) dt \right| \\
&\leq O_p(n^{-\frac{2-v}{2}-vm} + n^{-2vm-1/2}) + O(n^{-\frac{1-v}{2}-vm} + n^{-2vm}).
\end{aligned}$$

Lastly, it follows from the property of B-spline that $\|g_n^{(m)}(t)\|_{L^2} \leq b_3$, for some constant b_3

depending on $\|g_0^{(m)}(t)\|_{L^2}$ and m . Thus, we have

$$\begin{aligned} |I_5| &= |\langle W_\lambda g_n, h_n \rangle_\lambda| \\ &\leq \lambda \|g_n^{(m)}\|_{L^2} \|h_n\|_1 \\ &= \lambda O_p(n^{-\frac{1-v}{2}} + n^{-vm}) = o_p(n^{-(1-v)} + n^{-2vm}). \end{aligned}$$

As a result, we can conclude that $\alpha H'_n(\alpha) < 0$. Further, it is not hard to see that

$$H''_n(\alpha) = - \int_{\mathbb{I}} \exp(g_n(t) + \alpha h_n(t)) h_n^2(t) S_n(t) dt - \langle W_\lambda h_n, h_n \rangle_\lambda \leq 0,$$

which implies $H'_n(\alpha)$ is a nonincreasing function. Hence, $\hat{g}_{n,\lambda} \in [g_n - \alpha h_n, g_n + \alpha h_n]$. Note that

$$\begin{aligned} \|\hat{g}_{n,\lambda} - g_0\|_\infty &\leq \|\hat{g}_{n,\lambda} - g_n\|_\infty + \|\hat{g}_n - g_0\|_\infty \\ &\leq \alpha \|h_n\|_\infty + O(n^{-vm}) = O(n^{-\frac{1-v}{2}} + n^{-vm}), \end{aligned}$$

which goes to zero as $n \rightarrow \infty$. Recall that any two norms in the finite dimension Space are equivalent. Then, $\|\hat{g}_{n,\lambda} - g_n\|_1 = O(\|\hat{g}_{n,\lambda} - g_n\|_\infty) = O(n^{-\frac{1-v}{2}} + n^{-vm})$. Therefore, we have

$$\begin{aligned} \|\hat{g}_{n,\lambda} - g_0\|_1 &\leq \|\hat{g}_{n,\lambda} - g_n\|_1 + \|g_n - g_0\|_1 \\ &\leq O(n^{-\frac{1-v}{2}} + n^{-vm}) + \|g_n - g_0\|_1. \end{aligned}$$

By $\|g_n - g_0\|_\infty = O(n^{-vm})$, $\|g_n^{(m)}(t)\|_{L^2} \leq b_3$ and $g_0 \in \mathcal{H}^m$, we have

$$\|\hat{g}_{n,\lambda} - g_0\|_1 < \tilde{C},$$

and

$$J(\hat{g}_{n,\lambda} - g_0, \hat{g}_{n,\lambda} - g_0) < \tilde{C},$$

where \tilde{C} only depends on g_0 and m . The proof of Lemma 1 is complete.

Proof of Lemma 2. Following from equation (3.1) and Theorem 2 of Hoeffding (1963), we have

$$P(\|\mathcal{Z}_n(g) - \mathcal{Z}_n(f)\|_\lambda \geq t) \leq 2 \exp\left(\frac{-t^2}{8\|f - g\|_\infty}\right).$$

Together with Lemma 2.2.1 of Van Der Vaart and Wellner (1996), we have

$$\left\| \|\mathcal{Z}_n(f) - \mathcal{Z}_n(g)\|_\lambda \right\|_{\psi_2} \leq 8\|f - g\|_\infty,$$

where $\|\cdot\|_{\psi_2}$ denotes the orlicz norm associated with $\psi_2(s) = \exp(s^2) - 1$. Applying Theorem 2.2.4 of Van Der Vaart and Wellner (1996), for any $\varepsilon > 0$, we have

$$\begin{aligned} & \left\| \sup_{f, g \in \mathcal{G}, \|f - g\|_\infty \leq \varepsilon} \|\mathcal{Z}_n(g) - \mathcal{Z}_n(f)\|_\lambda \right\|_{\psi_2} \\ & \leq C' \left(\int_0^\varepsilon \sqrt{\log \{1 + N(\iota, \mathcal{G}, \|\cdot\|_\infty)\}} + \iota \sqrt{\log [1 + \{N(\iota, \mathcal{G}, \|\cdot\|_\infty)\}^2]} \right) d\iota \\ & \approx \varepsilon^{1 - \frac{1}{2m}}, \end{aligned}$$

where $N(\iota, \mathcal{G}, \|\cdot\|_\infty)$ is the covering number, the minimum number of $\|\cdot\|_\infty$ ι -balls needed to cover \mathcal{G} . Thus,

$$P \left(\sup_{g \in \mathcal{G}, \|g\|_\infty \leq \varepsilon} \|\mathcal{Z}_n(g)\|_\lambda \geq t \right) \leq 2 \exp(\varepsilon^{-2+1/m} t^2).$$

For brevity, we denote $\gamma \equiv 1 - 1/(2m)$, $T_n \equiv \{5 \log \log(n)\}^{1/2}$, $b_n = \sqrt{n}$, $\epsilon = b_n^{-1}$, and $Q_\epsilon = [-\log(\epsilon) - 1]$.

Note that

$$\begin{aligned}
& P\left(\sup_{g \in \mathcal{G}} \frac{b_n \|\mathcal{Z}_n(g)\|_\lambda}{b_n \|g\|_\infty^\gamma + 1} \geq T_n\right) \\
& \leq P\left(\sup_{g \in \mathcal{G}, \|g\|_\infty \leq \epsilon^{1/\gamma}} \frac{b_n \|\mathcal{Z}_n(g)\|_\lambda}{b_n \|g\|_\infty^\gamma + 1} \geq T_n\right) \\
& \quad + \sum_{l=0}^{Q_\epsilon} P\left(\sup_{g \in \mathcal{G}, (e^l \epsilon)^{1/\gamma} \leq \|g\|_\infty \leq (e^{l+1} \epsilon)^{1/\gamma}} \frac{b_n \|\mathcal{Z}_n(g)\|_\lambda}{b_n \|g\|_\infty^\gamma + 1} \geq T_n\right) \\
& \leq P\left(\sup_{g \in \mathcal{G}, \|g\|_\infty \leq \epsilon^{1/\gamma}} b_n \|\mathcal{Z}_n(g)\|_\lambda \geq T_n\right) \\
& \quad + \sum_{l=0}^{Q_\epsilon} P\left(\sup_{g \in \mathcal{G}, \|g\|_\infty \leq \{e^{l+1} \epsilon\}^{1/\gamma}} \frac{b_n \|\mathcal{Z}_n(g)\|_\lambda}{b_n \|g\|_\infty^\gamma + 1} \geq T_n\right) \\
& \leq 2 \exp\left\{-\left(\epsilon^{1/\gamma}\right)^{-2+1/m} T_n^2/n\right\} \\
& \quad + 2 \sum_{l=1}^{Q_\epsilon} \exp\left(-\left[\{e^{l+1} \epsilon\}^{1/\gamma}\right]^{-2+1/m} T_n^2 (e^l + 1)^2/n\right) \\
& = 2 \exp(-T_n^2) + 2 \sum_{l=1}^{Q_\epsilon} 2 \exp\{-e^{-2(l+1)T_n^2} (e^l + 1)^2\} \\
& \leq 2(Q_\epsilon + 2) \exp(-T_n^2/4) \\
& = b_5 \log(n) \{\log(n)\}^{-5/4} \xrightarrow{n \rightarrow \infty} 0,
\end{aligned}$$

where b_5 is a constant. We complete the proof of Lemma 2.

Proof of Theorem 1. Denote $g = \hat{g}_{n,\lambda} - g_0$. By Lemma 1, it is clear that $g \in \mathcal{G}$. Write

$$\begin{aligned}
l_{n,\lambda}(g_0 + g) - l_{n,\lambda}(g_0) &= \mathcal{S}_{n,\lambda}(g_0)g + \frac{1}{2}D\mathcal{S}_{n,\lambda}(g_0)gg + \frac{1}{6}D^2\mathcal{S}_{n,\lambda}(g^*)ggg \\
&\equiv I_1 + I_2 + I_3,
\end{aligned} \tag{S0.1}$$

where $g^* = g_0 + \alpha_1 g$ and $\alpha_1 \in [0, 1]$. We will next discuss the order of each term in (S0.1).

From the definition of $\hat{g}_{n,\lambda}$, we can get that $l_{n,\lambda}(g_0 + g) - l_{n,\lambda}(g_0) \geq 0$.

First, it follows from $g \in \mathcal{G}$ that there exists a constant \tilde{c} such that $\exp\{|g(t)|\} \leq \tilde{c}$. Then, we have

$$\begin{aligned}
 |6I_3| &= |D^2 \mathcal{S}_{n,\lambda}(g^*) g g g| \\
 &= \left| \frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t) + \alpha_1 g(t)\} g^3(t) I(Y_i \geq t) dt \right| \\
 &\leq \|g\|_\infty \left| \frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t) + \alpha_1 g(t)\} g^2(t) I(Y_i \geq t) dt \right| \\
 &\leq \tilde{c} \|g\|_\infty \left| \frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) I(Y_i \geq t) dt \right| \\
 &\leq \frac{\tilde{c} \|g\|_\infty}{n} \left| \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) I(Y_i \geq t) dt - n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S(t) dt \right| \\
 &\quad + \tilde{c} \|g\|_\infty \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S(t) dt. \tag{S0.2}
 \end{aligned}$$

Denote $\psi(Y; g) = \int_{\mathbb{I}} \exp\{g_0(t)\} g(t) I(Y \geq t) K_t dt$ and $\tilde{\psi}(Y; g) = b_2^{-1} c_m^{-1} h^{1/2} \psi(Y; g)$. Then, we have

$$\begin{aligned}
 &\left| \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) I(Y_i \geq t) dt - n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S(t) dt \right| \\
 &= \{b_2^{-1} c_m^{-1} h^{1/2}\}^{-1} \left| \sum_{i=1}^n \langle \tilde{\psi}(Y_i; g), g \rangle_\lambda - \langle E \tilde{\psi}(Y_i; g), g \rangle_\lambda \right|. \tag{S0.3}
 \end{aligned}$$

Under condition (C2), we have

$$\begin{aligned}
 \|\tilde{\psi}(Y; g) - \tilde{\psi}(Y; f)\|_\lambda &= b_2^{-1} c_m^{-1} h^{1/2} \|\psi(Y; g) - \psi(Y; f)\|_\lambda \\
 &= b_2^{-1} c_m^{-1} h^{1/2} \left\| \int_{\mathbb{I}} \exp\{g_0(t)\} \{f(t) - g(t)\} I(Y \geq t) K_t dt \right\|_\lambda \\
 &\leq b_2^{-1} c_m^{-1} h^{1/2} \int_{\mathbb{I}} \exp\{g_0(t)\} \|K_t\|_\lambda dt \|f - g\|_\infty \\
 &\leq b_2^{-1} c_m^{-1} h^{1/2} b_2 c_m h^{-1/2} \|f - g\|_\infty \\
 &= \|f - g\|_\infty.
 \end{aligned}$$

Together with Lemma 2, the following inequality holds with probability one

$$\left\| \sum_{i=1}^n \tilde{\psi}(Y_i; g) - E \tilde{\psi}(Y_i; g) \right\|_\lambda \leq \sqrt{n} \left\{ \|g\|_\infty^{1-1/(2m)} + 1 \right\} \{5 \log \log(n)\}^{1/2}. \tag{S0.4}$$

For the first term in equation (S0.1), it directly follows from (S0.2) and (S0.3) that

$$\begin{aligned}
& \left| \frac{\tilde{c}\|g\|_\infty}{n} \left| \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) I(Y_i \geq t) dt - n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S(t) dt \right| \right| \\
& \leq \frac{\tilde{c}\|g\|_\infty}{n} \{b_2^{-1} c_m^{-1} h^{1/2}\}^{-1} \|g\|_\lambda \left\| \sum_{i=1}^n \tilde{\psi}(Y_i; g) - E\tilde{\psi}(Y_i; g) \right\|_\lambda \\
& = O_p[\{5 \log \log(n)\}^{1/2} n^{-1/2} h^{-1}] \|g\|_\lambda^2.
\end{aligned}$$

As $\log\{\log(n)\}/nh^2 \rightarrow 0$, we have

$$\left| \frac{\tilde{c}\|g\|_\infty}{n} \left| \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) I(Y_i \geq t) dt - n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S(t) dt \right| \right| = o_p(1) \|g\|_\lambda^2.$$

For the second term in (S0.1), we have

$$\begin{aligned}
\tilde{c}\|g\|_\infty \left| \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S(t) dt \right| &= \tilde{c}\|g\|_\infty V(g, g) \\
&\leq \tilde{c}\|g\|_\infty \|g\|_\lambda^2.
\end{aligned}$$

Thus, we have $|6I_3| = o_p(1) \|g\|_\lambda^2$. It then follows from the Cauchy-Schwarz inequality that

$$|I_1| = |\mathcal{S}_{n,\lambda}(g_0)g| \leq \|\mathcal{S}_{n,\lambda}(g_0)\|_\lambda \|g\|_\lambda.$$

It follows from the proof of Corollary 1, for $\mathcal{S}_{n,\lambda}(g_0)$, we have

$$\begin{aligned}
\|\mathcal{S}_{n,\lambda}(g_0)\|_\lambda &= \left\| - \int_{\mathbb{I}} \exp\{g_0(t)\} K_t S_n(t) dt + \frac{1}{n} \sum_{i=1}^n \Delta_i K_{Y_i} - W_\lambda g_0 \right\|_\lambda \\
&\leq \frac{1}{n} \left\| \sum_{i=1}^n \int_{\mathbb{I}} K_t dM_i(t) \right\|_\lambda + \|W_\lambda g_0\|_\lambda = O_P((nh)^{-1/2} + \lambda^{1/2}).
\end{aligned}$$

Regarding I_2 , we have

$$\begin{aligned}
2I_2 &= D\mathcal{S}_{n,\lambda}(g_0)gg \\
&= D\mathcal{S}_{n,\lambda}(g_0)gg - D\mathcal{S}_\lambda(g_0)gg + D\mathcal{S}_\lambda(g_0)gg \\
&= -\|g\|_\lambda^2 + D\mathcal{S}_{n,\lambda}(g_0)gg - D\mathcal{S}_\lambda(g_0)gg.
\end{aligned}$$

This is because

$$\begin{aligned}
 & |DS_{n,\lambda}(g_0)gg - DS_\lambda(g_0)gg| \\
 = & \left| \frac{1}{n} \sum_{i=1}^n \left[\int_{\mathbb{I}} \exp\{g_0(t)\} I(Y_i \geq t) g^2(t) dt - \int_{\mathbb{I}} \exp\{g_0(t)\} S(t) g^2(t) dt \right] \right| \\
 \leq & \|g\|_\infty \left(\int_{\mathbb{I}} \exp\{g_0(t)\} < K_t, g >_\lambda \|S_n(t) - S(t)\| dt \right) \\
 = & \|g\|_\infty \|g\|_\lambda \left(\int_{\mathbb{I}} \exp\{g_0(t)\} \|K_t\|_\lambda \|S_n(t) - S(t)\|_\infty dt \right) \\
 = & o_p(1) O_p\{(nh)^{-1/2}\} \|g\|_\lambda.
 \end{aligned}$$

Therefore, we have

$$\|g\|_\lambda^2 \{1 + o_p(1)\} \leq \left[O_p\{(nh)^{-1/2} + \lambda^{1/2}\} + o_p\{(nh)^{-1/2}\} \right] \|g\|_\lambda,$$

which leads to $\|g\|_\lambda = O_p\{(nh)^{-1/2} + h^m\}$.

Proof of Theorem 2. For brevity, we denote $g = \hat{g}_{n,\lambda} - g_0$, $r_n = M\{(nh)^{-1/2} + h^m\}$, $\tilde{g} = d_n^{-1}g$ and $d_n = c_m r_n h^{-1/2}$. Recall $\|g\|_\lambda = O_p\{(nh)^{-1/2} + h^m\}$ in Theorem 1. Then, there exists a constant M such that the event $B_n = \{\|g\|_\lambda \leq r_n\}$ happens with large probability. Since $h = o(1)$ and $\log\{\log(n)\}/(nh^2) \rightarrow 0$ as $n \rightarrow \infty$, it is easy to see that $d_n = o(1)$. On the other hand, when event B_n happens, one can get $\|\tilde{g}\|_\infty \leq 1$ and

$$J(\tilde{g}, \tilde{g}) = d_n^{-2} \lambda^{-1} \lambda J(g, g) \leq d_n^{-2} \lambda^{-1} \|g\|_\lambda^2 \leq d_n^{-2} \lambda^{-1} r_n^2 \leq c_m^{-2} h \lambda^{-1}.$$

It then follows directly that $\tilde{g} \in \mathcal{F}$, where $\mathcal{F} = \{g : \|g\|_\infty \leq 1, J(g, g) \leq c_m^{-2} h \lambda^{-1}\}$ when the

event B_n happens. Next, by a Taylor expansion, we have

$$\begin{aligned}
& \mathcal{S}_n(\hat{g}_{n,\lambda}) - \mathcal{S}_n(g_0) - \{\mathcal{S}(\hat{g}_{n,\lambda}) - \mathcal{S}(g_0)\} \\
&= -\frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t) + g(t)\} I(Y_i \geq t) K_t dt + \frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} I(Y_i \geq t) K_t dt \\
&\quad - \left[-\int_{\mathbb{I}} \exp\{g_0(t) + g(t)\} S(t) K_t dt + \int_{\mathbb{I}} \exp\{g_0(t)\} S(t) K_t dt \right] \\
&= -\frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} [\exp\{g(t)\} - 1] I(Y_i \geq t) K_t dt \\
&\quad \quad \quad + \int_{\mathbb{I}} \exp\{g_0(t)\} [\exp\{g(t)\} - 1] S(t) K_t dt \\
&= -\frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} \left[g(t) + \frac{g(t)^2}{2} \{1 + o_p(1)\} \right] \{I(Y_i \geq t) - S(t)\} K_t dt \\
&= -\frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g(t) \{I(Y_i \geq t) - S(t)\} K_t dt \\
&\quad - \frac{1}{2n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g(t)^2 \{I(Y_i \geq t) - S(t)\} K_t dt \{1 + o_p(1)\} \\
&\equiv I_1 + I_2.
\end{aligned}$$

Observe that

$$\begin{aligned}
\|2I_2\|_\lambda &= \left\| -\frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) \{I(Y_i \geq t) - S(t)\} K_t dt \{1 + o_p(1)\} \right\|_\lambda \\
&= \left\| \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) \{S_n(t) - S(t)\} K_t dt \{1 + o_p(1)\} \right\|_\lambda \\
&\leq \int_{\mathbb{I}} \exp\{g_0(t)\} \|K_t\|_\lambda dt \|S_n(t) - S(t)\|_\infty \|g\|_\infty^2 \\
&\leq O_p(n^{-1/2}) (c_m h^{-1/2})^2 \|g\|_\lambda^2 c_m h^{-1/2} \\
&= O_p\{(n^{1/2}h)^{-1}\} \|g\|_\lambda^2 c_m h^{-1/2}.
\end{aligned}$$

The fact $\log\{\log(n)\}/(nh^2) \rightarrow 0$ implies $nh^2 \rightarrow \infty$ as $n \rightarrow \infty$. Thus,

$$I_2 = o_p(1) c_m h^{-1/2} \|g\|_\lambda^2 = o_p(1) c_m h^{-1/2} \{(nh)^{-1/2} + h^m\}^2. \quad (\text{S0.5})$$

We next consider I_1 . Write

$$\begin{aligned}
-I_1 &= \frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g(t) \{I(Y_i \geq t) - S(t)\} K_t dt \\
&= \frac{1}{n} \sum_{i=1}^n \phi(Y_i, g) - E\{\phi(Y_i, g)\},
\end{aligned}$$

where $\phi(Y, g) = \int_{\bar{1}} \exp\{g_0(t)\}g(t)I(Y \geq t)K_t dt$. We denote

$\tilde{\phi}(Y; \tilde{g}) = b_2^{-1}(c_m h^{-1/2})^{-1} d_n^{-1} \phi(Y, d_n \tilde{g})$. By a careful evaluation, we can show that

$$\|\tilde{\phi}(Y; \tilde{g}) - \tilde{\phi}(Y; \tilde{f})\|_\lambda \leq \|\tilde{f} - \tilde{g}\|_\infty.$$

Define

$$\mathcal{Z}_n(g) \equiv \frac{1}{\sqrt{n}} \sum_{i=1}^n \phi(Y_i, g) - E\phi(Y_i, g).$$

In light of Lemma S.1 in Shang and Cheng (2013) and the proof of Lemma 2, we have

$$\lim_{n \rightarrow \infty} P \left[\sup_{g \in \mathcal{F}} \frac{\|\mathcal{Z}_n(g)\|_\lambda}{h^{-(2m-1)/4m} \|g\|_\infty^{1-1/(2m)} + n^{-1/2}} \leq \{5 \log \log(n)\}^{1/2} \right] = 1.$$

Then, we have

$$\begin{aligned} \|I_1\|_\lambda &= b_2(c_m h^{-1/2})d_n \left\| \frac{1}{n} \sum_{i=1}^n \tilde{\phi}(Y_i, \tilde{g}) - E\{\tilde{\phi}(Y_i, \tilde{g})\} \right\|_\lambda \\ &= \frac{b_2(c_m h^{-1/2})d_n}{n} \left\{ \sqrt{n} \|\tilde{g}\|_\infty^{1-1/(2m)} h^{-(2m-1)/(4m)} + 1 \right\} \{5 \log \log(n)\}^{1/2}. \end{aligned}$$

In view of the fact that $m > 1/2$ and $\|\tilde{g}\|_\infty \leq 1$, we get

$$\begin{aligned} & \frac{b_2(c_m h^{-1/2})d_n}{n} \left\{ \sqrt{n} \|\tilde{g}\|_\infty^{1-1/(2m)} h^{-(2m-1)/(4m)} + 1 \right\} \{5 \log \log(n)\}^{1/2} \\ &= O_p(h^{-(6m-1)/(4m)} \{n^{-\frac{1}{2}} + h^{(2m-1)/(4m)}\} \{5 \log \log(n)\}^{\frac{1}{2}} \{(nh)^{-1/2} + h^m\}) \\ &= O_p(h^{-\frac{6m-1}{4m}} n^{-1/2} \{\log \log(n)\}^{1/2} \{(nh)^{-1/2} + h^m\}). \end{aligned} \tag{S0.6}$$

Hence, combing (S0.5) and (S0.6), we have

$$\begin{aligned} & \mathcal{S}_n(\hat{g}_{n,\lambda}) - \mathcal{S}_n(g_0) - \{\mathcal{S}(\hat{g}_{n,\lambda}) - \mathcal{S}(g_0)\} \\ &= O_p(h^{-\frac{6m-1}{4m}} n^{-\frac{1}{2}} \{\log \log(n)\}^{\frac{1}{2}} \{(nh)^{-\frac{1}{2}} + h^m\} + h^{-\frac{1}{2}} \{(nh)^{-1} + h^{2m}\}). \end{aligned}$$

On the other hand, we note that

$$\begin{aligned}
& \mathcal{S}_n(\hat{g}_{n,\lambda}) - \mathcal{S}_n(g_0) - \{\mathcal{S}(\hat{g}_{n,\lambda}) - \mathcal{S}(g_0)\} \\
&= \mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda}) - \mathcal{S}_{n,\lambda}(g_0) - \{\mathcal{S}_\lambda(\hat{g}_{n,\lambda}) - \mathcal{S}_\lambda(g_0)\} \\
&= \mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda}) - \mathcal{S}_{n,\lambda}(g_0) - \left\{ D\mathcal{S}_\lambda(g_0)g + \int_{\mathbb{I}} \int_{\mathbb{I}} sD^2\mathcal{S}_\lambda(g_0 + s'sg)g^2 ds ds' \right\} \\
&= g - \mathcal{S}_{n,\lambda}(g_0) - \int_{\mathbb{I}} \int_{\mathbb{I}} sD^2\mathcal{S}_\lambda(g_0 + s'sg)g^2 ds ds' + \mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda}).
\end{aligned}$$

For any $h \in \mathcal{H}^m$, there exists $h_n \in \Psi_{m,\mathcal{I}}$ such that $\|h - h_n\|_\infty = O(n^{-vm})$. Furthermore, by the definition of $\mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda})$, we have $\mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda})h_n = 0$. Then, we further write

$$\begin{aligned}
& \mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda})h \\
&= \mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda})(h - h_n) \\
&= \mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda})(h - h_n) - \mathcal{S}_n(g_0)(h - h_n) + \mathcal{S}_n(g_0)(h - h_n) \\
&= \left[- \int_{\mathbb{I}} \exp\{g_0(t)\}g(t)\{1 + o_p(1)\}\{h(t) - h_n(t)\}S_n(t) dt - \langle W_\lambda g_0, h - h_n \rangle_\lambda \right] \\
&\quad - \int_{\mathbb{I}} \exp\{g_0(t)\}\{h(t) - h_n(t)\}S_n(t) dt + \frac{1}{n} \sum_{i=1}^n \Delta_i(h - h_n)(Y_i) - \langle W_\lambda g_0, h - h_n \rangle_\lambda \\
&\equiv L_1 + L_2 + L_3.
\end{aligned}$$

First, we consider L_1 . Write

$$\begin{aligned}
|L_1| &= \left| \int_{\mathbb{I}} \exp\{g_0(t)\}g(t)\{1 + o_p(1)\}\{h(t) - h_n(t)\}S_n(t) dt + \langle W_\lambda g_0, h - h_n \rangle_\lambda \right| \\
&\leq \left| \left[\int_{\mathbb{I}} \exp\{g_0(t)\}g^2(t)S_n(t) dt \right]^{1/2} \left[\int_{\mathbb{I}} \exp\{g_0(t)\}\{h(t) - h_n(t)\}^2 S_n(t) dt \right]^{1/2} \right| \\
&\quad \times \{1 + o_p(1)\} + |\langle W_\lambda g_0, h - h_n \rangle_\lambda| \\
&\leq O_p(\|g\|_\lambda \|h - h_n\|_\infty) + o_p(\lambda n^{-vm}) \\
&= O_p(n^{-vm} \{(nh)^{-1/2} + h^m\}) + o_p(h^{2m} n^{-vm}).
\end{aligned}$$

Second, we consider L_2 and obtain

$$\begin{aligned}
 E\{|L_2|^2\} &= E\left[\left|\int_{\mathbb{I}} \exp\{g_0(t)\}\{h(t) - h_n(t)\}S_n(t) dt - \frac{1}{n} \sum_{i=1}^n \Delta_i\{h(Y_i) - h_n(Y_i)\}\right|^2\right] \\
 &= \frac{1}{n} E\left[\left|\int_{\mathbb{I}} \exp\{g_0(t)\}\{h(t) - h_n(t)\} dM_i(t)\right|^2\right] \\
 &\leq \frac{b_4}{n} \|h - h_n\|_{\infty}^2 \\
 &= O(n^{-1-2vm}),
 \end{aligned}$$

where b_4 is a constant. which implies $|L_2| = O_p(n^{-1/2-vm})$. Third, it is not hard to verify that

$|L_3| = O_p(\lambda n^{-vm})$. Combining the asymptotic order of L_1, L_2 and L_3 , we have

$$\|\mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda})\|_{\lambda} = O_p(n^{-1/2-vm} + h^{2m} n^{-vm} + n^{-vm}\{(nh)^{-1/2} + h^m\}).$$

Lastly, we observe

$$\left\|\int_{\mathbb{I}} \int_{\mathbb{I}} s D^2 \mathcal{S}_{\lambda}(g_0 + s' s g) g^2 ds ds'\right\|_{\lambda} \leq \int_{\mathbb{I}} \int_{\mathbb{I}} \|D^2 \mathcal{S}_{\lambda}(g_0 + s' s g) g^2\|_{\lambda} ds ds'.$$

In particular,

$$\begin{aligned}
 \|D^2 \mathcal{S}_{\lambda}(g_0 + s' s g) g^2\|_{\lambda} &= \left\|\int_{\mathbb{I}} \exp\{g_0(t) + s s' g(t)\} S(t) g^2(t) K_t dt\right\|_{\lambda} \\
 &\leq \tilde{c}(c_m h^{-1/2}) \|g\|_{\lambda}^2.
 \end{aligned}$$

Finally, we get

$$\|g - \mathcal{S}_{n,\lambda}(g_0)\|_{\lambda} = O_p(\alpha_n),$$

where

$$\begin{aligned}
 \alpha_n &= n^{-1/2-vm} + n^{-vm}((nh)^{-1/2} + h^m) + h^{-1/2}((nh)^{-1} + h^{2m}) \\
 &\quad + h^{-(6m-1)/(4m)} n^{-1/2} (\log \log(n))^{1/2} ((nh)^{-1/2} + h^m).
 \end{aligned}$$

The proof of Theorem 2 is complete.

Proof of Theorem 3. For ease of presentation, we denote $R_n = \hat{g}_{n,\lambda} - g^* - \mathcal{S}_n(g_0)$ and $g^* = (id - W_{\lambda})g_0$. It follows directly from Theorem 2 that $\|R_n\|_{\lambda} = O_p(\alpha_n) = o_p(n^{-1/2})$. It

can be checked that $\|\mathcal{S}_n\|_\lambda = O_p((nh)^{-1/2})$. Hence, R_n is asymptotically negligible compared with \mathcal{S}_n . In the following, we derive the asymptotic distribution of $(nh)^{-1/2}\{\hat{g}_{n,\lambda}(t_0) - g^*(t_0)\}$. Recall that for any $t \in \mathbb{I}$ and $g \in \mathcal{H}^m$, we have $\langle K_t, g \rangle_\lambda = g(t)$. Then, we have

$$\begin{aligned} |(nh)^{1/2} \langle K_{t_0}, \hat{g}_{n,\lambda} - g^* - \mathcal{S}_n(g_0) \rangle_\lambda| &\leq \|K_{t_0}\|_\lambda \|R_n\|_\lambda (nh)^{1/2} \\ &\leq c_m h^{-1/2} (nh)^{1/2} o_p(n^{-1/2}) \\ &= o_p(1). \end{aligned}$$

Next, we write

$$\begin{aligned} -(nh)^{1/2} \langle K_{t_0}, \mathcal{S}_n(g_0) \rangle_\lambda &= (nh)^{1/2} \left[\int_{\mathbb{I}} \exp\{g_0(t)\} \mathcal{S}_n(t) K_t(t_0) dt - \frac{1}{n} \sum_{i=1}^n \Delta_i K_{Y_i}(t_0) \right] \\ &= (nh)^{1/2} \left[\int_{\mathbb{I}} \exp\{g_0(t)\} \mathcal{S}_n(t) K_{t_0}(t) dt - \frac{1}{n} \sum_{i=1}^n \Delta_i K_{t_0}(Y_i) \right] \\ &= (nh)^{1/2} \frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} K_{t_0}(t) dM_i(t) \\ &= \frac{1}{\sqrt{nh^{-1}}} \sum_{i=1}^n \int_{\mathbb{I}} K_{t_0}(t) dM_i(t). \end{aligned}$$

Observe that

$$\text{Var} \left\{ \int_{\mathbb{I}} K_{t_0}(t) dM_i(t) \right\} = \int_{\mathbb{I}} K_{t_0}^2(t) \exp\{g_0(t)\} S(t) dt = V(K_{t_0}, K_{t_0}).$$

Invoking $hV(K_{t_0}, K_{t_0}) < h\|K\|_\lambda^2 < c_m^2$ and $hV(K_{t_0}, K_{t_0}) \rightarrow \sigma_{t_0}^2$ as $n \rightarrow \infty$, we have

$$(nh)^{1/2} \langle K_{t_0}, \mathcal{S}_n(g_0) \rangle_\lambda \xrightarrow{d} N(0, \sigma_{t_0}^2)$$

as $n \rightarrow \infty$. The proof of Theorem 3 is complete.

Proof of Corollary 1. First, for any t ,

$$W_\lambda g_0(t) = \langle W_\lambda g_0, K_t \rangle_\lambda = \sum_{j=0}^{\infty} \frac{\lambda \gamma_j}{1 + \lambda \gamma_j} h_j(t) V(g_0, h_j).$$

By the Cauchy-Schwarz's inequality,

$$\begin{aligned} \sum_{j=0}^{\infty} \frac{\lambda\gamma_j}{1+\lambda\gamma_j} h_j(t) V(g_0, h_j) &\leq \left\{ \sum_{j=0}^{\infty} \lambda\gamma_j V^2(g_0, h_j) \right\}^{\frac{1}{2}} \left\{ \sum_{j=0}^{\infty} \frac{\lambda\gamma_j}{(1+\lambda\gamma_j)^2} h_j^2(t) \right\}^{\frac{1}{2}} \\ &\leq h^m \sup_{j \in N} \|h_j\|_{\infty} \sqrt{J(g_0, g_0)} \left\{ \sum_{j=0}^{\infty} \frac{\lambda\gamma_j}{(1+\lambda\gamma_j)^2} \right\}^{\frac{1}{2}}. \end{aligned}$$

Invoking $g_0 \in \mathcal{H}^m$ and $\gamma_j \asymp j^{2m}$, $W_{\lambda}g_0(t) = O(h^{m-1/2})$. Hence,

$$\sqrt{nh}W_{\lambda}g_0(t) = O(n^{1/2}h^m) = o(1).$$

It follows directly from Theorem 3 that the results of Corollary 1 hold.

Proof of Theorem 5 (ii). For notational convenience, we denote $\hat{g} = \hat{g}_{n,\lambda}$, $\hat{g}^0 = \hat{g}_{n,\lambda}^0$, $g = \hat{g}^0 + \omega_0 - \hat{g}$. By Theorem 4,

$$\|g\|_{\lambda} = \|\hat{g}^0 + \omega_0 - \hat{g}\|_{\lambda} \leq \|\hat{g}^0 + \omega_0 - g_0\|_{\lambda} + \|\hat{g} - g_0\|_{\lambda} = O_p(r_n),$$

where $r_n = (nh)^{-1/2} + h^m$. Applying Taylor expansion,

$$\begin{aligned} \text{LRT}_{n,\lambda} &= L_{n,\lambda}(\omega_0 + \hat{g}^0) - \mathbb{L}_{n,\lambda}(\hat{g}) \\ &= \mathcal{S}_{n,\lambda}(\hat{g})(\omega_0 + \hat{g}^0 - \hat{g}) + \int_{\mathbb{I}} \int_{\mathbb{I}} s D\mathcal{S}_{n,\lambda}(\hat{g} + ss'g) gg \, ds \, ds'. \end{aligned}$$

It follows from the definition of $\mathcal{S}_{n,\lambda}(\hat{g})$ that $\mathcal{S}_{n,\lambda}(\hat{g})(\hat{g}^0 + \omega_0 - \hat{g}) = 0$. Hence, we write

$$\begin{aligned} \text{LRT}_{n,\lambda} &= \int_{\mathbb{I}} \int_{\mathbb{I}} s D\mathcal{S}_{n,\lambda}(\hat{g} + ss'g) gg \, ds \, ds' \\ &= \int_{\mathbb{I}} \int_{\mathbb{I}} s \{ D\mathcal{S}_{n,\lambda}(\hat{g} + ss'g) gg - D\mathcal{S}_{n,\lambda}(g_0) gg \} \, ds \, ds' + \int_{\mathbb{I}} \int_{\mathbb{I}} s D\mathcal{S}_{n,\lambda}(g_0) gg \, ds \, ds' \\ &= \int_{\mathbb{I}} \int_{\mathbb{I}} s \{ D\mathcal{S}_{n,\lambda}(\hat{g} + ss'g) gg - D\mathcal{S}_{n,\lambda}(g_0) gg \} \, ds \, ds' \\ &\quad + \frac{1}{2} \{ D\mathcal{S}_{n,\lambda}(g_0) gg - D\mathcal{S}_{\lambda}(g_0) gg \} + \frac{1}{2} D\mathcal{S}_{\lambda}(g_0) gg \\ &\equiv I_1 + I_2 + I_3. \end{aligned}$$

We first consider I_1 . Denote $\tilde{g} = \hat{g} + ss'g - g_0$ for any $0 \leq s, s' \leq 1$. Then, $\|\tilde{g}\|_{\lambda} = O_p(r_n)$ and

$$\begin{aligned} D\mathcal{S}_{n,\lambda}(\hat{g} + ss'g) gg &= D\mathcal{S}_{n,\lambda}(\tilde{g} + g_0) gg \\ &= - \int_{\mathbb{I}} \exp\{g_0(t) + \tilde{g}(t)\} g(t) g(t) S_n(t) \, dt - \langle W_{\lambda}g, g \rangle_{\lambda}. \end{aligned}$$

Thus, we have

$$\begin{aligned}
& |DS_{n,\lambda}(\hat{g} + ss'g)gg - DS_{n,\lambda}(g_0)gg| \\
&= \left| - \int_{\mathbb{I}} [\exp\{g_0(t) + \hat{g}(t)\} - \exp\{g_0(t)\}] g(t)g(t)S_n(t) dt \right| \\
&= \left| \int_{\mathbb{I}} \exp\{g_0(t)\} \tilde{g}(t) \{1 + o_p(1)\} g^2(t) S_n(t) dt \right| \\
&\leq \|\tilde{g}\|_{\infty} \left\| \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S_n(t) dt \right\| \\
&\leq \|\tilde{g}\|_{\infty} \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) |S_n(t) - S(t)| dt + \|\tilde{g}\|_{\infty} \left| \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S(t) dt \right|.
\end{aligned}$$

Under condition (C3) and the assumption that $nh^4 \rightarrow \infty$ as $n \rightarrow \infty$, we have

$$\begin{aligned}
& \|\tilde{g}\|_{\infty} \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) |S_n(t) - S(t)| dt \\
&= \|\tilde{g}\|_{\infty} \int_{\mathbb{I}} \exp\{g_0(t)\} g(t) \langle K_t, g \rangle_{\lambda} |S_n(t) - S(t)| dt \\
&\leq \|\tilde{g}\|_{\infty} \|K_t\|_{\lambda} \|g\|_{\lambda} \|S_n(t) - S(t)\|_{\infty} \|g(t)\|_{\infty} \int_{\mathbb{I}} \exp\{g_0(t)\} dt \\
&= O_p(n^{-1/2}h^{-1}) \|g\|_{\lambda}^2 \|\tilde{g}\|_{\infty} \\
&= o_p(1) \|g\|_{\lambda}^2 \|\tilde{g}\|_{\infty}.
\end{aligned}$$

Moreover, note that

$$\|\tilde{g}\|_{\infty} \left\| \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) S(t) dt \right\| \leq \|\tilde{g}\|_{\infty} \|g\|_{\lambda}^2,$$

which gives that $|I_1| = O_p(1) \|\tilde{g}\|_{\infty} \|g\|_{\lambda}^2 = O_p(h^{-1/2}r_n^3)$.

We next consider I_2 . Write

$$\begin{aligned}
2|I_2| &= |DS_{n,\lambda}(g_0)gg - DS_{\lambda}(g_0)gg| \\
&= \frac{1}{n} \left| \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g^2(t) \{I(Y_i \geq t) - S(t)\} dt \right| \\
&= \frac{1}{n} \left\langle \sum_{i=1}^n \int_{\mathbb{I}} \exp\{g_0(t)\} g(t) K_t \{I(Y_i \geq t) - S(t)\} dt, g \right\rangle_{\lambda}.
\end{aligned}$$

We can show along the same lines of Theorem 2 that $|I_2| = O_p(r_n a'_n)$ and

$$a'_n = h^{-(6m-1)/(4m)} n^{-1/2} \{\log \log(n)\}^{1/2} \{(nh)^{-1/2} + h^m\}.$$

Lastly, we consider I_3 . Applying the fact that $I_3 = -\|g\|_\lambda^2/2$ and combining the previous arguments, we have

$$\text{LRT}_{n,\lambda} = -\frac{\|g\|_\lambda^2}{2} + O_p(h^{-1/2}r_n^3 + r_n a'_n).$$

Recall that $nh^{2m} \rightarrow 0$. Hence, $nh^{2m+1} \rightarrow 0$ as $n \rightarrow \infty$. Together with $nh^4 \rightarrow \infty$, we have $h^{-1/2}r_n^3 + r_n a'_n = o(n^{-1})$. As a result,

$$-2n\text{LRT}_{n,\lambda} = n\|\hat{g}^0 + \omega_0 - \hat{g}\|_\lambda^2 + o_p(1).$$

Proof of Theorem 5 (iii). In view of $-2n\text{LRT}_{n,\lambda} = n\|\hat{g}^0 + \omega_0 - \hat{g}\|_\lambda^2 + o_p(1)$ in part (ii) of Theorem 5, to show part (iii), it suffices to derive the asymptotic properties of $n\|\hat{g}^0 + \omega_0 - \hat{g}\|_\lambda^2$.

It is not hard to see that

$$\begin{aligned} & n^{1/2} \|\hat{g}^0 + \omega_0 - \hat{g} - \mathcal{S}_{n,\lambda}^0(g_0^0) + \mathcal{S}_{n,\lambda}(g_0)\|_\lambda \\ & \leq n^{1/2} \|\hat{g}^0 + \omega_0 - \mathcal{S}_{n,\lambda}^0(g_0^0)\|_\lambda + n^{1/2} \|\hat{g} - \mathcal{S}_{n,\lambda}(g_0)\|_\lambda \\ & = O_p(n^{1/2}a_n) = o_p(1). \end{aligned}$$

Thus, we only need to focus on $n^{1/2}\{\mathcal{S}_{n,\lambda}^0(g_0^0) - \mathcal{S}_{n,\lambda}(g_0)\}$. Recall that

$$\begin{aligned} \mathcal{S}_{n,\lambda}^0(g_0^0) &= -\int_{\mathbb{I}} \exp\{g_0(t)\} S_n(t) K_t^* dt + \frac{1}{n} \sum_{i=1}^n \Delta_i K_{Y_i}^* - W_\lambda^* g_0^0 \\ &= -\int_{\mathbb{I}} \exp\{g_0(t)\} S_n(t) \left\{ K_t - \frac{K_{t_0}(t)K_{t_0}}{K(t_0, t_0)} \right\} dt \\ &\quad + \frac{1}{n} \sum_{i=1}^n \Delta_i \left\{ K_{Y_i} - \frac{K_{t_0}(Y_i)K_{t_0}}{K(t_0, t_0)} \right\} - \left\{ W_\lambda g_0 - \frac{(W_\lambda g_0)(t_0)}{K(t_0, t_0)} K_{t_0} \right\}. \end{aligned}$$

Then, we have

$$\mathcal{S}_{n,\lambda}^0(g_0^0) - \mathcal{S}_{n,\lambda}(g_0) = \frac{K_{t_0}}{K(t_0, t_0)} \left[\frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} K_{t_0}(t) dM_i(t) + (W_\lambda g_0)(t_0) \right],$$

and

$$\sqrt{n} \|\mathcal{S}_{n,\lambda}^0(g_0^0) - \mathcal{S}_{n,\lambda}(g_0)\|_\lambda = \left| \frac{1}{\sqrt{K(t_0, t_0)}} \left[\frac{1}{\sqrt{n}} \sum_{i=1}^n \int_{\mathbb{I}} K_{t_0}(t) dM_i(t) + \sqrt{n}(W_\lambda g_0)(t_0) \right] \right|.$$

Applying $nh^{2m} \rightarrow 0$, we get

$$\frac{\sqrt{n}(W_\lambda g_0)(t_0)}{\|K_{t_0}\|_\lambda} \leq \frac{\sqrt{nh}(W_\lambda g_0)(t_0)}{h^{1/2}\|V^{1/2}(K_{t_0}, K_{t_0})\|_\lambda} = O(1) \frac{\sqrt{nh}(W_\lambda g_0)(t_0)}{\sigma_{t_0}} = O(\sqrt{nh^m}) = o(1).$$

Combining these gives

$$\frac{1}{\sqrt{K(t_0, t_0)}} \left\{ \frac{1}{\sqrt{n}} \sum_{i=1}^n \int_{\mathbb{I}} K_{t_0}(t) dM_i(t) + \sqrt{n}(W_\lambda g_0)(t_0) \right\} \xrightarrow{d} N(0, c_{t_0}),$$

where

$$c_{t_0} = \lim_{h \rightarrow 0} \frac{V(K_{t_0}, K_{t_0})}{\|K_{t_0}\|^2} \in (0, 1].$$

As a result, it follows immediately that $-2n\text{LRT}_{n,\lambda} \xrightarrow{d} c_{t_0}\chi_1^2$ implying $\|\hat{g}^0 + \omega_0 - \hat{g}\|_\lambda = O_p(n^{-1/2})$. Thereby, we prove the first part of Theorem 5. The proof of Theorem 5 is complete.

Proof of Theorem 6. For simplicity, we denote $g = g_0 - \hat{g}_{n,\lambda}$ and $r_n = (nh)^{-1/2} + h^m$. By a Taylor expansion, we get

$$\begin{aligned} \text{PLRT}_{n,\lambda} &= l_{n,\lambda}(g_0) - l_{n,\lambda}(\hat{g}_{n,\lambda}) \\ &= \mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda})(g_0 - \hat{g}_{n,\lambda}) + \int_{\mathbb{I}} \int_{\mathbb{I}} s D\mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda} + ss'g) ds ds' \\ &\equiv I_1 + I_2. \end{aligned}$$

We first consider I_1 . Along similar lines of the proof of Theorem 2, we have

$$\begin{aligned} |I_1| &= |\mathcal{S}_{n,\lambda}(\hat{g}_{n,\lambda})g| \\ &\leq \|\mathcal{S}_{n,\lambda}\|_\lambda \|g\|_\lambda \\ &= O_p[n^{-1/2-vm} + h^{2m}v^{-vm} + n^{-vm}\{(nh)^{-1/2} + h^m\}]\|g\|_\lambda \\ &= O_p(r_n[n^{-1/2-vm} + n^{-vm}\{(nh)^{-1/2} + h^m\}]). \end{aligned}$$

Similar to the proof of Theorem 5(ii), it can be easily verified that

$$|I_2| = -\frac{\|g\|_\lambda^2}{2} + O_p(h^{-1/2}r_n^3 + r_n\alpha'_n),$$

where $\alpha'_n = h^{-(6m-1)/(4m)} n^{-1/2} \{\log \log(n)\}^{1/2} r_n$. Thus,

$$\text{PLRT}_{n,\lambda} = -\frac{\|g\|_\lambda^2}{2} + O_p(h^{-1/2} r_n^3 + r_n \alpha''_n),$$

where $\alpha''_n = \alpha'_n + n^{-1/2-vm} + n^{-vm} \{(nh)^{-1/2} + h^m\}$. Under the conditions that $m > (3 + \sqrt{5})/4$, $1/(4m) \leq v \leq 1/(2m)$, $nh^{2m+1} = O(1)$ and $nh^3 \rightarrow \infty$, we have

$$-2n\text{PLRT}_{n,\lambda} = n\|g\|_\lambda^2 + o_p(h^{-1/2}).$$

On the other hand, under H_0^{global} , g_0 is true function. Then, by Theorems 2 and 3, we have $\|\hat{g}_{n,\lambda} - g_0 - \mathcal{S}_{n,\lambda}(g_0)\| = O_p(\alpha_n)$ and $n^{1/2}\alpha_n = o(1)$. Combining these gives

$$n^{1/2}\|g\|_\lambda = n^{1/2}\|\mathcal{S}_{n,\lambda}(g_0)\|_\lambda + o_p(1).$$

Next, we consider $\|\mathcal{S}_{n,\lambda}(g_0)\|_\lambda$. Through direct calculations,

$$n\|\mathcal{S}_{n,\lambda}(g_0)\|_\lambda^2 = n^{-1} \left\| \sum_{i=1}^n \int_{\mathbb{I}} K_t dM_i(t) \right\|_\lambda^2 + 2 < \sum_{i=1}^n \int_{\mathbb{I}} K_t dM_i(t), W_{\lambda g_0} > + n\|W_{\lambda g_0}\|_\lambda^2.$$

We first approximate $\|W_{\lambda g_0}\|_\lambda$. To this end, we define $m_\lambda(j) \equiv |V(g_0, h_j)|^2 \gamma_j \frac{\lambda \gamma_j}{1 + \lambda \gamma_j}$ and $m(j) \equiv |V(g_0, h_j)|^2 \gamma_j$, $j = 0, 1, 2, \dots$. Note that $|m_\lambda(j)|$ is a sequence of functions satisfying $|m_\lambda(j)| \leq m(j)$. Since $g_0 \in \mathcal{H}^m$,

$$\sum_j |V(g_0, h_j)|^2 \gamma_j = \int_N m(j) d\mu(j) = J(g_0, g_0) < \infty,$$

where $\mu(\cdot)$ is the counting measure. Invoking $\lim_{\lambda \rightarrow 0} m_\lambda(j) = 0$,

$$\sum_j |V(g_0, h_j)|^2 \frac{\lambda \gamma_j^2}{1 + \lambda \gamma_j} = \int_N m_\lambda(j) d\mu(j) \rightarrow 0$$

as $\lambda \rightarrow 0$ by the Lebesgue dominated convergence theorem. That is,

$$\|W_{\lambda g_0}\|_\lambda^2 = \sum_j |V(g_0, h_j)|^2 \frac{\lambda^2 \gamma_j^2}{1 + \lambda \gamma_j} = o(\lambda).$$

Using this fact, we have

$$\begin{aligned} E \left| < \sum_{i=1}^n \int_{\mathbb{I}} K_t dM_i(t), W_{\lambda g_0} > \right|^2 &= E \left| \sum_{i=1}^n \int_{\mathbb{I}} W_{\lambda g_0}(t) dM_i(t) \right|^2 \\ &= n \int_{\mathbb{I}} \exp\{g_0(t)\} S(t) \{W_{\lambda}(g_0(t))\}^2 dt \\ &\leq n\|W_{\lambda}(g_0(t))\|_\lambda^2 = o(n\lambda). \end{aligned}$$

Together with $nh^{2m+1} = O(1)$, it follows that

$$\left\langle \sum_{i=1}^n \int_{\mathbb{I}} K_t dM_i(t), W_\lambda g_0 \right\rangle = o_p\{(n\lambda)^{1/2}\} = o_p(n^{1/2}h^m) = o_p(h^{-1/2}).$$

So far, we have shown that

$$n\|\mathcal{S}_{n,\lambda}(g_0)\|_\lambda^2 = n^{-1} \left\| \sum_{i=1}^n \int_{\mathbb{I}} K_t dM_i(t) \right\|_\lambda^2 + o_p(h^{-1}).$$

In what follows, we derive the limiting distribution of $n^{-1} \left\| \sum_{i=1}^n \int_{\mathbb{I}} K_t dM_i(t) \right\|_\lambda^2$. A direct calculation yields that

$$\frac{1}{n} \left\| \sum_{i=1}^n \int_{\mathbb{I}} K_t dM_i(t) \right\|_\lambda^2 = \frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_i(s) + \frac{1}{n} W_n,$$

where $W_n = \sum_{i \neq j} \int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_j(s)$. Denoting $W_{ij} = 2 \int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_j(s)$, one can write $W_n = \sum_{1 \leq i < j \leq n} W_{ij}$. So, W_n is clean [Jong (1987)]. Next, we aim to derive the limiting distribution of W_n . Let $\sigma_n^2 = Var(W_n)$. Write

$$\begin{aligned} \sigma_n^2 &= \frac{n(n-1)}{2} E(W_{ij}^2) \\ &= 2n(n-1) E \left\{ \int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_j(s) \right\}^2 \\ &= 2n(n-1) \sum_{l=0}^{\infty} \frac{1}{(1+\lambda\gamma_l)^2}. \end{aligned}$$

More notation is needed here. Define G_1, G_2 and G_4 as follows:

$$G_1 \equiv \sum_{i < j} E(W_{ij}^4),$$

$$G_2 \equiv \sum_{i < j < k} \{E(W_{ij}^2 W_{ik}^2) + E(W_{ji}^2 W_{jk}^2) + E(W_{ki}^2 W_{kj}^2)\},$$

$$G_3 \equiv \sum_{i < j < k < l} \{E(W_{ij} W_{ik} W_{lj} W_{lk}) + E(W_{ij} W_{il} W_{kj} W_{kl}) + E(W_{ik} W_{il} W_{jk} W_{jl})\}.$$

By Proposition 3.2 of Jong (1987), if G_1, G_2, G_3 are all of lower order than σ_n^4 , $\sigma_n^{-1} W_n$ converges weakly to the standard normal distribution. Now, we study the order of each $G_i, i = 1, 2, 3$.

First, observe that

$$\begin{aligned}
& E\{W_{ij}^4\} \\
&= 16E\left\{\int_{\mathbb{I}}\int_{\mathbb{I}}\langle K_t, K_s \rangle dM_i(t) dM_j(s)\right\}^4 \\
&= 16\int_{\mathbb{I}}\int_{\mathbb{I}}\int_{\mathbb{I}}\int_{\mathbb{I}}\int_{\mathbb{I}}\int_{\mathbb{I}}\int_{\mathbb{I}}\int_{\mathbb{I}}\langle K_{t_1}, K_{s_1} \rangle \langle K_{t_2}, K_{s_2} \rangle \langle K_{t_3}, K_{s_3} \rangle \langle K_{t_4}, K_{s_4} \rangle \\
&\quad E\left\{dM_i(t_1) dM_j(s_1) dM_i(t_2) dM_j(s_2) dM_i(t_3) dM_j(s_3) dM_i(t_4) dM_j(s_4)\right\} \\
&\leq 16\|K_{t_1}\|_{\lambda}^4 \|K_{s_1}\|_{\lambda}^4 \int_{\mathbb{I}}[E\{dM_i(t_1)\}]^4 \int_{\mathbb{I}}[E\{dM_j(s_1)\}]^4 \\
&= O(h^{-4}),
\end{aligned}$$

which implies $G_1 = O(n^2 h^{-4})$. Next, by Cauchy-Schwarz inequity,

$$E\{W_{ij}^2 W_{ik}^2\} \leq [E\{W_{ij}^4\}]^{1/2} [E\{W_{ik}^4\}]^{1/2} = O(h^{-4}),$$

which gives $G_2 = O(n^3 h^{-4})$. A straightforward calculation yields that

$$E(W_{ij} W_{ik} W_{lj} W_{lk}) = 16 \sum_{j=0}^{\infty} \frac{1}{(1 + \lambda \gamma_j)^4} = O(h^{-1}).$$

Therefore, $G_3 = O(n^4 h^{-1})$. Combining the fact that $\sigma_n^4 = (\sigma_n^2)^2 = O(n^4 h^{-2})$ and the assumptions that $nh^3 \rightarrow \infty$ and $h = o(1)$, G_1, G_2, G_3 are of lower order than that of σ_n^4 . Hence, by Jong (1987),

$$\sigma_n^{-1} W_n \xrightarrow{d} N(0, 1)$$

as $n \rightarrow \infty$. Recall that $\rho_{\lambda}^2 = \sum_{j=0}^{\infty} \frac{h}{(1 + \lambda \gamma_j)^2}$. We have

$$\frac{1}{\sqrt{2h^{-1}n\rho_{\lambda}}} W_n \xrightarrow{d} N(0, 1). \quad (\text{S0.7})$$

Lastly, we consider $\frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_i(s)$. By a direct calculation,

$$E\left\{\int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_i(s)\right\}^2 = O(\|K_t\|_{\lambda}^4) = O(h^{-2}).$$

Then,

$$\begin{aligned}
& E\left\{\sum_{i=1}^n \int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_i(s) - h^{-1} \sigma_{\lambda}^2\right\}^2 \\
&\leq nE\left\{\int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_i(s)\right\}^2 = O(nh^{-2}),
\end{aligned}$$

where $\sigma_\lambda^2 = \sum_{j=0}^{\infty} \frac{h}{1+\lambda\gamma_j}$. Combining these gives

$$\frac{1}{n} \sum_{i=1}^n \int_{\mathbb{I}} \int_{\mathbb{I}} \langle K_t, K_s \rangle dM_i(t) dM_i(s) = h^{-1} \sigma_\lambda^2 + O_p\{(n^{1/2}h)^{-1}\}. \quad (\text{S0.8})$$

Combining (S0.7) and (S0.8), we have $n\|\mathcal{S}_{n,\lambda}\|_\lambda^2 = O_p(h^{-1})$ and therefore $n^{1/2}\|\mathcal{S}_{n,\lambda}\|_\lambda = O_p(h^{-1/2})$. As a result,

$$\begin{aligned} -2n\text{PLRT}_{n,\lambda} &= \{n^{1/2}\|\mathcal{S}_{n,\lambda}\|_\lambda + o_p(1)\}^2 + o_p(h^{-1/2}) \\ &= n\|\mathcal{S}_{n,\lambda}\|_\lambda^2 + o_p(h^{-1/2}). \end{aligned} \quad (\text{S0.9})$$

In view of (S0.7), (S0.8) and (S0.9), we have that as $n \rightarrow \infty$,

$$(2h^{-1}\sigma_\lambda^4/\rho_\lambda^2)^{-1/2} \{-2n\gamma_\lambda\text{PLRT}_{n,\lambda} - n\gamma_\lambda\|W_\lambda g_0(t)\|_\lambda^2 - h^{-1}\sigma_\lambda^4/\rho_\lambda^2\} \xrightarrow{d} N(0, 1).$$

Proof of Theorem 7. First, it can be easily verified that $m > (3 + \sqrt{5})/4$, $1/(4m) \leq v \leq 1/(2m)$ and $h \asymp n^{-d}$ with $1/(2m+1) \leq d < 1/3$ satisfy those conditions in Theorem 6. Throughout this proof, we only consider $g_{n_0} = g_0 + g_n$ for $g_n \in \mathcal{A}$ in H_1 . To prove Theorem 7, we write

$$\begin{aligned} -2n \cdot \text{PLRT}_{n,\lambda} &= -2n\{l_{n,\lambda}(g_0) - l_{n,\lambda}(g_{n_0})\} - 2n\{l_{n,\lambda}(g_{n_0}) - l_{n,\lambda}(\hat{g}_{n,\lambda})\} \\ &\equiv I_1 + I_2. \end{aligned} \quad (\text{S0.10})$$

We first consider I_1 . For simplicity, we denote

$$\begin{aligned} R_i &= \left[- \int_{\mathbb{I}} \exp\{g_0(t)\} I(Y_i \geq t) dt + \Delta_i g_0(Y_i) \right] - \\ &\quad \left[- \int_{\mathbb{I}} \exp\{g_{n_0}(t)\} I(Y_i \geq t) dt + \Delta_i g_{n_0}(Y_i) \right] \\ &= - \int_{\mathbb{I}} g_n(t) dM_i(t) - \int_{\mathbb{I}} \int_{\mathbb{I}} \exp\{g_{n_0}(t) - sg_n(t)\} g_n^2(t) I(Y_i \geq t) dt ds. \end{aligned}$$

Then,

$$\begin{aligned} E\{R_i^2\} &\leq 2 \int_{\mathbb{I}} g_n^2(t) S(t) \exp\{g_{n_0}(t)\} dt \\ &\quad + 2E \left[\int_{\mathbb{I}} \int_{\mathbb{I}} \exp\{g_{n_0}(t) - sg_n(t)\} g_n^2(t) I(Y_i \geq t) dt ds \right]^2 \\ &= O(\|g_n\|_{\lambda}^2 + \|g_n\|_{\lambda}^4). \end{aligned}$$

Therefore, we can get

$$E \left\{ \left| \sum_{i=1}^n (R_i - ER_i) \right|^2 \right\} \leq nE\{R_i^2\} = (n\|g_n\|_{\lambda}^2 + n\|g_n\|_{\lambda}^4).$$

Combining these gives

$$n [l_{n,\lambda}(g_0) - l_{n,\lambda}(g_{n_0}) - E\{l_{n,\lambda}(g_0) - l_{n,\lambda}(g_{n_0})\}] = O_p(n^{1/2}\|g_n\|_{\lambda} + n^{1/2}\|g_n\|_{\lambda}^2).$$

On the other hand, invoking $DS_{\lambda}(g)g_n g_n < 0$ for any $g \in \mathcal{H}^m$, there exists constant $c' > 0$ such that

$$\begin{aligned} E\{DS_{n,\lambda}(g_{n_0}^*)g_n g_n\} &\leq c' E\{DS_{n,\lambda}(g_{n_0})g_n g_n\} \\ &= -c' \|g_n\|_{\lambda}^2. \end{aligned}$$

Then, we have

$$\begin{aligned} E\{l_{n,\lambda}(g_0) - l_{n,\lambda}(g_{n_0})\} &= E \left\{ \mathcal{S}_{n,\lambda}(g_{n_0})(-g_n) + \frac{1}{2} DS_{n,\lambda}(g_{n_0}^*)g_n g_n \right\} \\ &\leq \lambda J(g_{n_0}, g_n) - \frac{c' \|g_n\|_{\lambda}^2}{2} \\ &\leq \{J(g_n, g_n) + J(g_0, g_n)\} - \frac{c' \|g_n\|_{\lambda}^2}{2} \\ &\leq \{J(g_n, g_n) + J(g_0, g_0)^{1/2} J(g_n, g_n)^{1/2}\} - \frac{c' \|g_n\|_{\lambda}^2}{2} \\ &= O(\lambda) - \frac{c' \|g_n\|_{\lambda}^2}{2}. \end{aligned}$$

It then follows that

$$\begin{aligned} I_1 &\geq n\|g_n\|_{\lambda}^2 + O_p(n\lambda + n^{1/2}\|g_n\|_{\lambda} + n^{1/2}\|g_n\|_{\lambda}^2) \\ &= n\|g_n\|_{\lambda}^2 \{1 + O_p(\lambda\|g_n\|_{\lambda}^{-2} + n^{-1/2}\|g_n\|_{\lambda}^{-1} + n^{-1/2})\}. \end{aligned} \tag{S0.11}$$

Second, we consider I_2 . Under alternative hypothesis H_{1n} , note that $\|\hat{g}_{n,\lambda} - g_{n_0}\| = O_p\{(nh)^{-1/2} + h^m\}$. It then follows by the FBR in Theorem 2 that

$$\inf_{n \geq N} \inf_{g_n \in \mathcal{A}} P_{g_{n_0}} (\|\hat{g}_{n,\lambda} - g_{n_0} - S_{n,\lambda}(g_{n_0})\|_\lambda \leq Mr_n) \rightarrow 1, \quad (\text{S0.12})$$

where $r_n = (nh)^{-1/2} + h^m$, $P_{g_{n_0}}$ means the probability relies on g_{n_0} . Along the lines of Theorem 6, we can show I_2 has the same limiting distribution as in Theorem 6, uniformly for any $g_n \in \mathcal{A}$. In other words, uniformly over all $g_n \in \mathcal{A}$,

$$(2\nu_{n_0})^{-1/2}(I_2 - n\|W_\lambda g_{n_0}\|_\lambda^2 - h^{-1}\sigma_{n_0,\lambda}^2) = O_p(1), \quad (\text{S0.13})$$

where $\nu_{n_0} = h^{-1}\sigma_{n_0,\lambda}^4/\rho_{n_0,\lambda}^2$, $\sigma_{n_0,\lambda}^2$ and $\rho_{n_0,\lambda}^2$ are defined the same as σ_λ^2 and ρ_λ^2 but with eigenvalues and eigenvectors obtained under g_{n_0} . Next, let V_{n_0} and V_0 be functions defined similarly to V in Section 2. Thus, for any $f \in \mathcal{H}^m$,

$$\begin{aligned} |V_{n_0}(f, f) - V_0(f, f)| &= \left| \int_{\mathbb{I}} [\exp\{g_{n_0}(t)\} - \exp\{g_0(t)\}] S(t) f^2(t) dt \right| \\ &\leq \|\exp\{g_n(t)\}\|_\infty V_0(f, f) \|g_n\|_\infty \\ &= \zeta V_0(f, f) \|g_n\|_\infty. \end{aligned}$$

It follows from the supplementary material (page 56) of Shang and Cheng (2013) that

$$\sigma_{n_0,\lambda}^2 - \sigma_\lambda^2 = O(h^{-1/2}\|g_n\|_\lambda). \quad (\text{S0.14})$$

Combining (S0.11), (S0.13) and (S0.14) gives

$$\begin{aligned} &(2\nu_n)^{-1/2}(-2nr_\lambda \text{PLRT}_{n,\lambda} - \nu_n) \\ &= (2\nu_n)^{-1/2}\{-r_\lambda(I_1 + I_2) - \nu_n\} \\ &= (2\nu_n)^{-1/2}r_\lambda(I_2 - n\|W_\lambda g_{n_0}\|_\lambda^2 - h^{-1}\sigma_{n_0,\lambda}^2) + (2\nu_n)^{-1/2}r_\lambda n\|W_\lambda g_{n_0}\|_\lambda^2 \\ &\quad + (2\nu_n)^{-1/2}r_\lambda I_1 + (2\nu_n)^{-1/2}r_\lambda h^{-1}(\sigma_{n_0,\lambda}^2 - \sigma_\lambda^2) \\ &\geq O_p(1) + (2\nu_n)^{-1/2}r_\lambda n\|g_n\|_\lambda^2 \{1 + O_p(\lambda\|g_n\|_\lambda^{-2} + n^{-1/2}\|g_n\|_\lambda^{-1} + n^{-1/2})\} \\ &\quad + O(h^{-1}\|g_n\|_\lambda), \end{aligned}$$

where $O_p(1)$ holds uniformly in \mathcal{A} , $\nu_n = h^{-1}\sigma_\lambda^4/\rho_\lambda^2$ and r_λ is defined in Theorem 6. Let $\lambda\|g_n\|_\lambda^{-2} \leq 1/C$, $n^{-1/2}\|g_n\|_\lambda^{-1} \leq 1/C$, $Ch^{-1}\|g_n\|_\lambda \leq (nh^{1/2})\|g_n\|_\lambda^2$, and $\|g_n\|_\lambda^2 \geq C(nh^{1/2})^{-1}$ for some sufficiently small constant C . In other words,

$$|(2\nu_n)^{-1/2}(-2nr_\lambda\text{PLRT}_{n,\lambda} - \nu_n)| \geq c_\alpha,$$

where c_α is the critical value (based on $N(0,1)$) to H_0^{global} at nominal level α . This leads to

$$\|g_n\|_\lambda^2 \geq C\{h^{2m} + (nh^{1/2})^{-1}\}. \quad (\text{S0.15})$$

Combining (S0.12) and (S0.15), we complete the proof of Theorem 7.